

**"MAGNETOHYDRODYNAMIC
MODELING OF CORONAL EVOLUTION
AND DISRUPTION"**

NASW-01005

3rd Quarter (1st Year) Progress Report,

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MAGNETOHYDRODYNAMIC MODELING OF CORONAL EVOLUTION AND DISRUPTION

3rd QUARTER (1st YEAR) PROGRESS REPORT:

3/30/02 - 6/29/02

Our progress for the 3rd quarter of 2002 was summarized in an invited review talk on CMEs given by Jon Linker at the Solar Wind X conference. Slides from the talk are attached in the following pages.

MODELING CMES IN THE CORONA AND SOLAR WIND*

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*Supported by NASA and NSF

OUTLINE

- 1) Introduction
- 2) Overview of CME initiation models
 - a) Types of models
 - b) Storage and release models
 - c) What we do / do not understand
- 3) Overview of interplanetary CME models
 - a) Types of models
 - b) Can interplanetary observations distinguish between CME initiation models?
 - c) How well do we understand interplanetary flux ropes?
- 4) Summary: My guess at the Future

INTRODUCTION

- Coronal Mass Ejections (CMEs) are a fundamental aspect of solar and heliospheric physics.
- Despite many years of study, their origin and evolution is poorly understood:
 - We don't know how CMEs are initiated in the corona.
 - We don't know how they give rise to the structures we observe in interplanetary space.
- Present observations, as well as new observations that will be available in the next few years, give us the opportunity to make significant progress on these problems. Modeling is a key ingredient to success.
- This will be a “narrow” review: I will try to touch on the areas that I believe represent the upcoming challenges in modeling CMEs, and where I think significant advances are likely to be made.

CLASSIFICATION OF MODELS

(Klimchuk 2001)

- **Storage and Release**
 - Energy is stored in the magnetic field over a long period of time (days to weeks), and released as a result of instability, loss of equilibrium, or nonequilibrium (cf., Forbes, JGR, 2001)
- **Directly Driven**
 - Energy is pumped into the corona during eruption
 - A flux rope structure is assumed: can be used to fit white light observations
 - No observational support for vast energy flux into corona at eruption (cf., Forbes, Spring AGU 2001)
- **Thermal Blast**
 - Thermal energy is input in the form of an unspecified energy source (*e.g.*, thermal energy from a flare)
 - Lots of observational problems, currently not in favor

STORAGE AND RELEASE MODELS

- Energy is stored over a long period and released over a short period
- Instability is a competition between magnetic field tension and magnetic pressure:
- For example, for force-free equilibria:

$$\mathbf{J} \times \mathbf{B} = 0$$

$$(\nabla \times \mathbf{B}) \times \mathbf{B} = 0$$

$$\mathbf{B} \cdot \nabla \mathbf{B} = \frac{1}{2} \nabla B^2$$

- Generally, eruption occurs when field line tension is reduced or when pressure is increased
- There must be free energy \Leftrightarrow parallel electric current \Leftrightarrow twist \Leftrightarrow shear
- Highly nonpotential magnetic structures are in fact frequently observed

HOW IS THE ENERGY STORED?

- Photospheric motions can store energy in the field by twisting / shearing.
- Magnetic fields may emerge already twisted (*i.e.*, carrying current) from below the photosphere.
- Recent studies (e.g., Demoulin et al., 2002) indicate that the twist in the field primarily emerges with new fields.
- Differential rotation is unlikely to provide the primary energization of the field; smaller scale motions are not yet ruled out.

STORAGE AND RELEASE MODELS: EXAMPLES

- Flux Cancellation Model (*e.g.*, van Ballegooijen & Martens 1989; Forbes & Isenberg 1991; Amari *et al.* 2000; Linker *et al.* 2001)
- Breakout Model (Antiochos, DeVore, & Klimchuk, 1999)
- A new model by Zhang and Low postulates that the rough classification of two types of CMEs (fast and slow) are related to “normal” and “inverse” polarity prominences
- It is difficult to distinguish between the models:
 - CME initiation does not produce significant photospheric magnetic field changes
 - In many models, the eruption is a threshold effect (% flux change, critical shear, etc.)
 - Differences between models can be very subtle

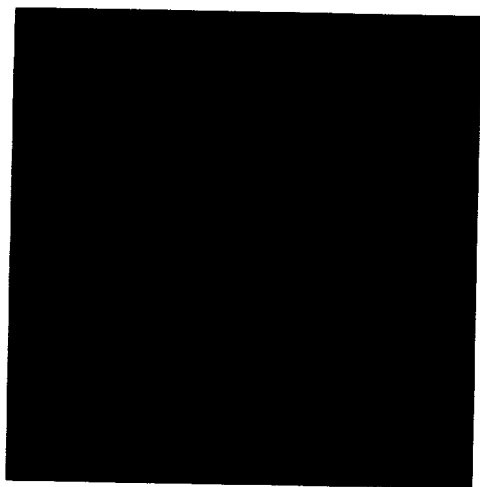
FLUX CANCELLATION MODEL

(e.g., van Ballegooijen & Martens 1989;
Forbes & Isenberg 1991;
Amari *et al.* 2000;
Linker *et al.* 2001)

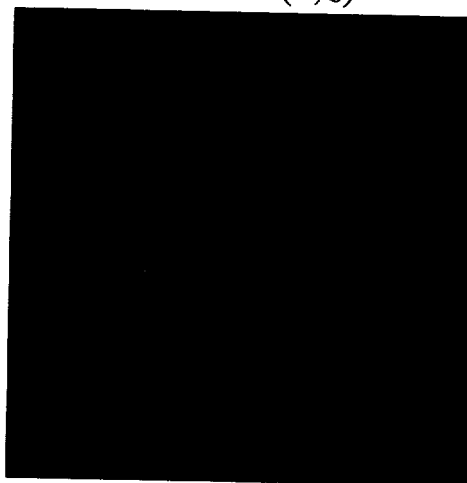
- Flux cancellation at the neutral line can destabilize a sheared arcade
- Flows that converge toward the neutral line can lead to flux cancellation (van Ballegooijen & Martens 1989)
- A flux rope forms above the neutral line
- The dips in the magnetic field lines can support prominence material
- This mechanism produces an energetic eruption with significant conversion of stored magnetic energy into kinetic energy
- There is a threshold for eruption: emergence of less flux than the threshold leads to the formation of a stable filament
- Even a small amount of emerged flux can trigger an eruption
- Dispersal of the magnetic flux in an active region can provide the necessary trigger

Eruption of a Helmet Streamer By Emerging Flux

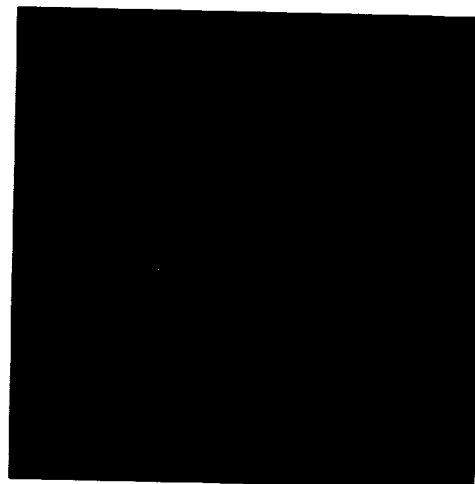
Flux $\Psi(r,z)$



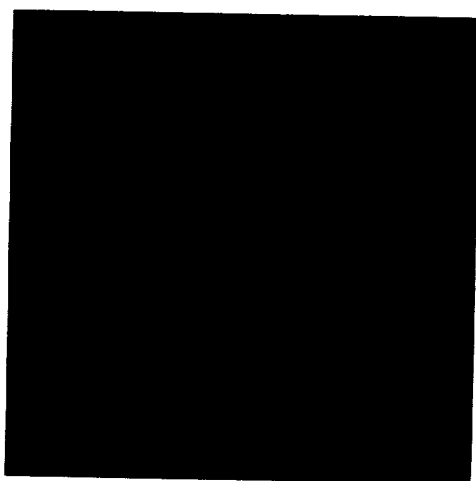
Unsheared streamer



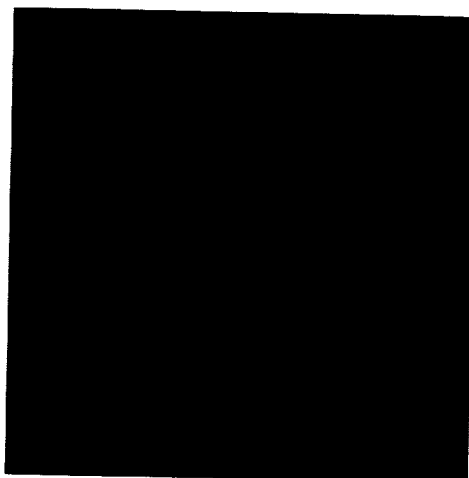
Sheared streamer
 $t = t_0$



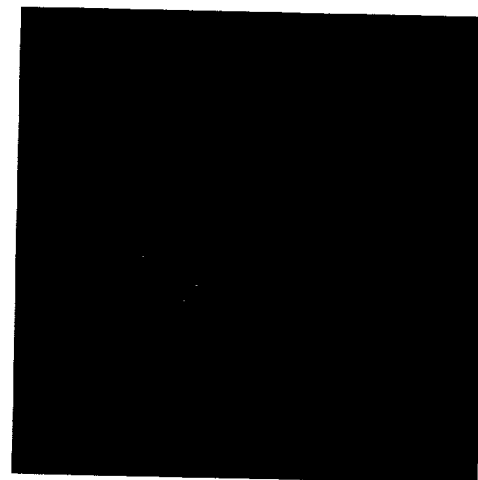
4.5% emerged flux
 $t = t_0 + 6$ hours



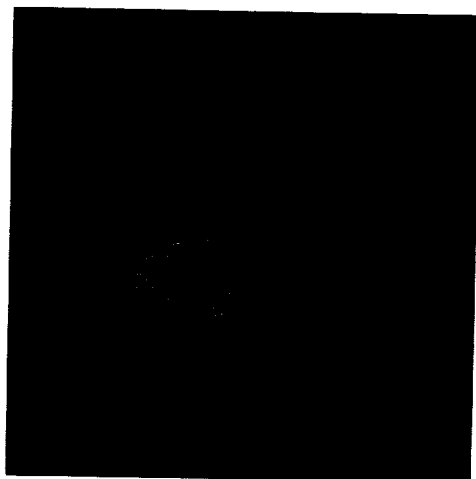
7.5% emerged flux
 $t = t_0 + 10$ hours



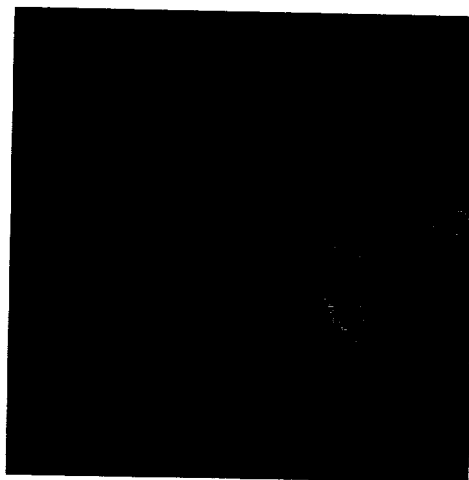
10.5% emerged flux
 $t = t_0 + 14$ hours



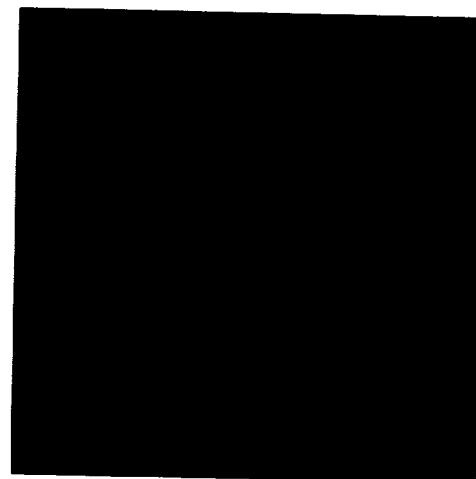
12% emerged flux
 $t = t_0 + 16$ hours



13.5% emerged flux
 $t = t_0 + 18$ hours



15% emerged flux
 $t = t_0 + 20$ hours

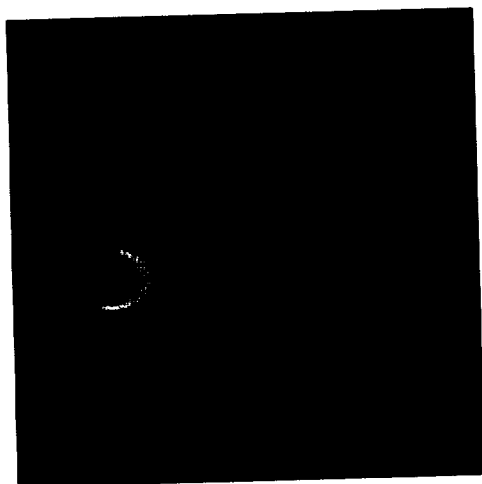


15% emerged flux
 $t = t_0 + 2.5$ days

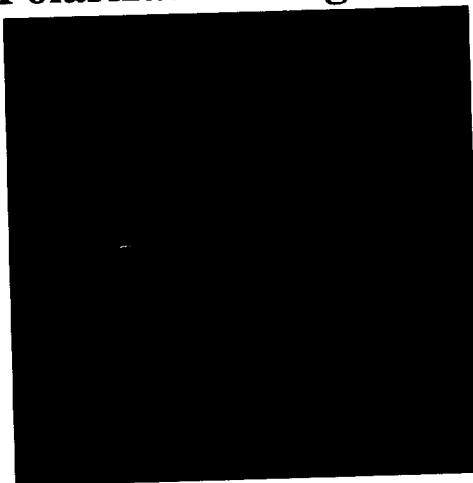
Eruption of a Helmet Streamer By Emerging Flux



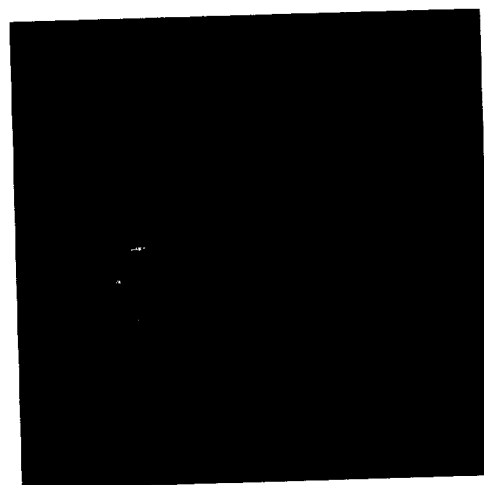
Polarization Brightness



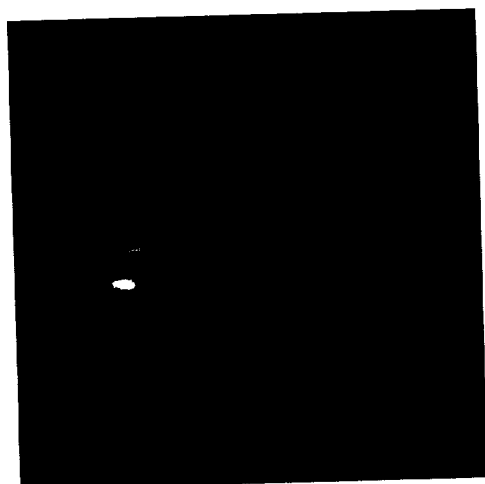
Unsheared streamer



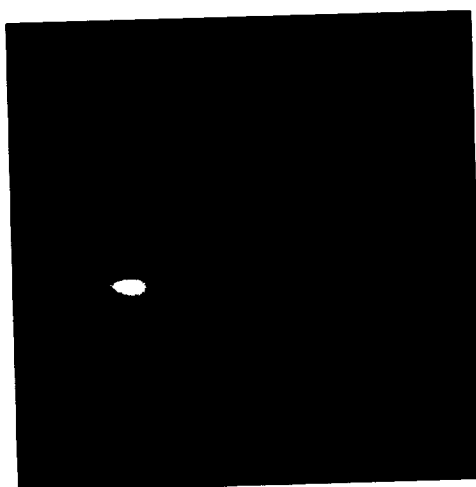
Sheared streamer
 $t = t_0$



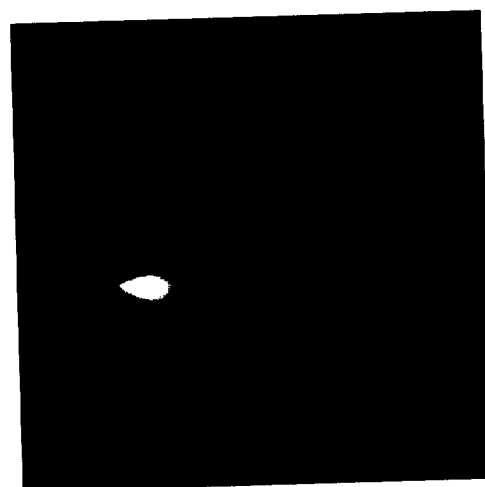
4.5% emerged flux
 $t = t_0 + 6$ hours



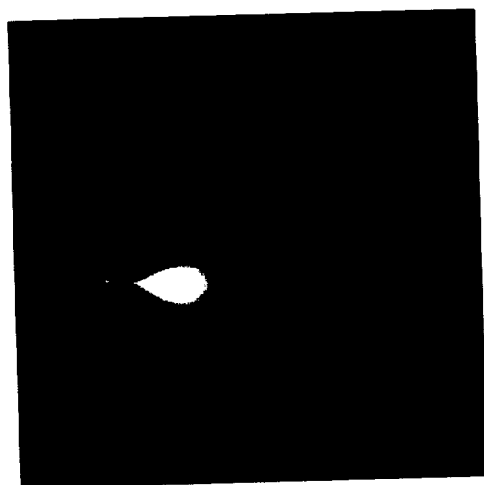
7.5% emerged flux
 $t = t_0 + 10$ hours



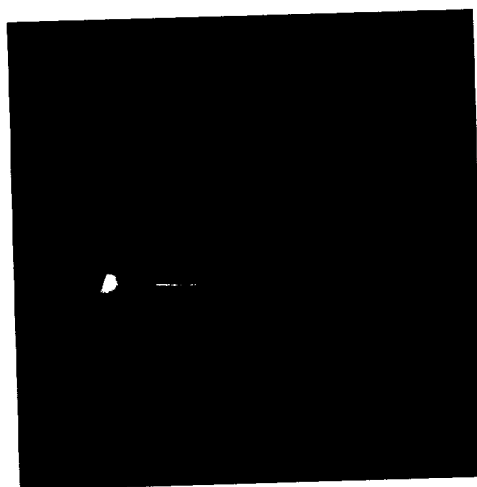
10.5% emerged flux
 $t = t_0 + 14$ hours



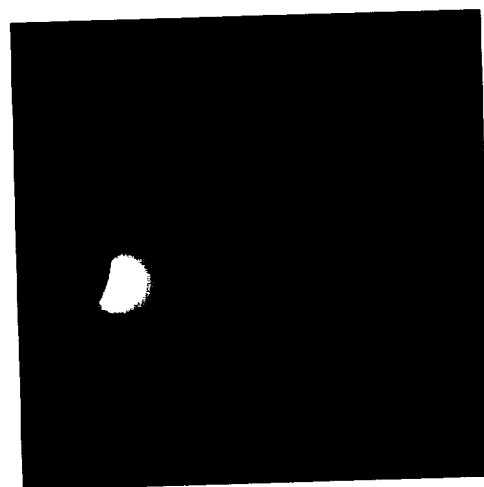
12% emerged flux
 $t = t_0 + 16$ hours



13.5% emerged flux
 $t = t_0 + 18$ hours



15% emerged flux
 $t = t_0 + 20$ hours



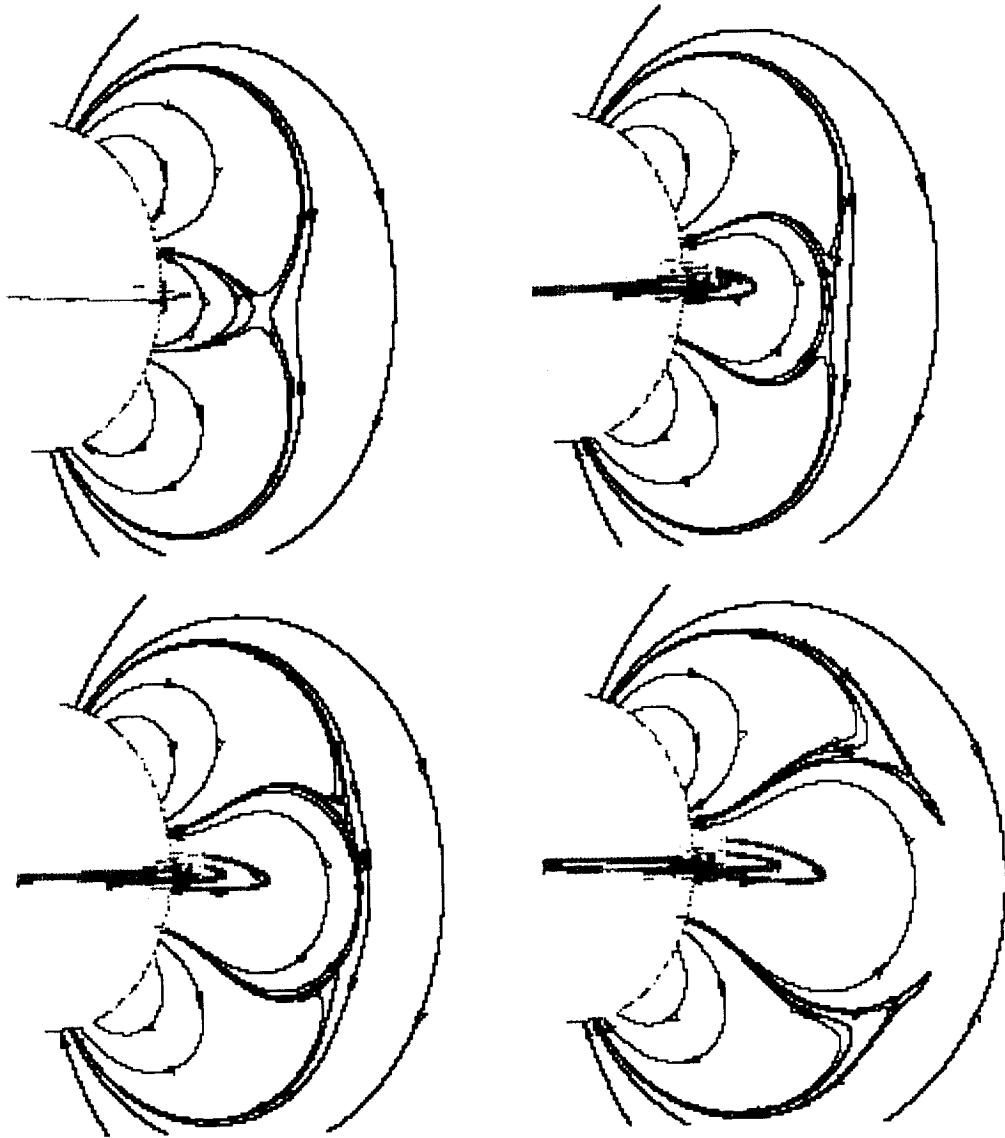
15% emerged flux
 $t = t_0 + 2.5$ days

BREAKOUT MODEL

(Antiochos, DeVore, & Klimchuk 1999,
Ap. J., 510, 485.)

- Requires a more complex magnetic field topology than a simple bipolar magnetic field
- Driven by increasing shear near the neutral line
- Eruption occurs when overlying magnetic field lines reconnect at an X-point, releasing the downward tension force

The "Breakout" Model

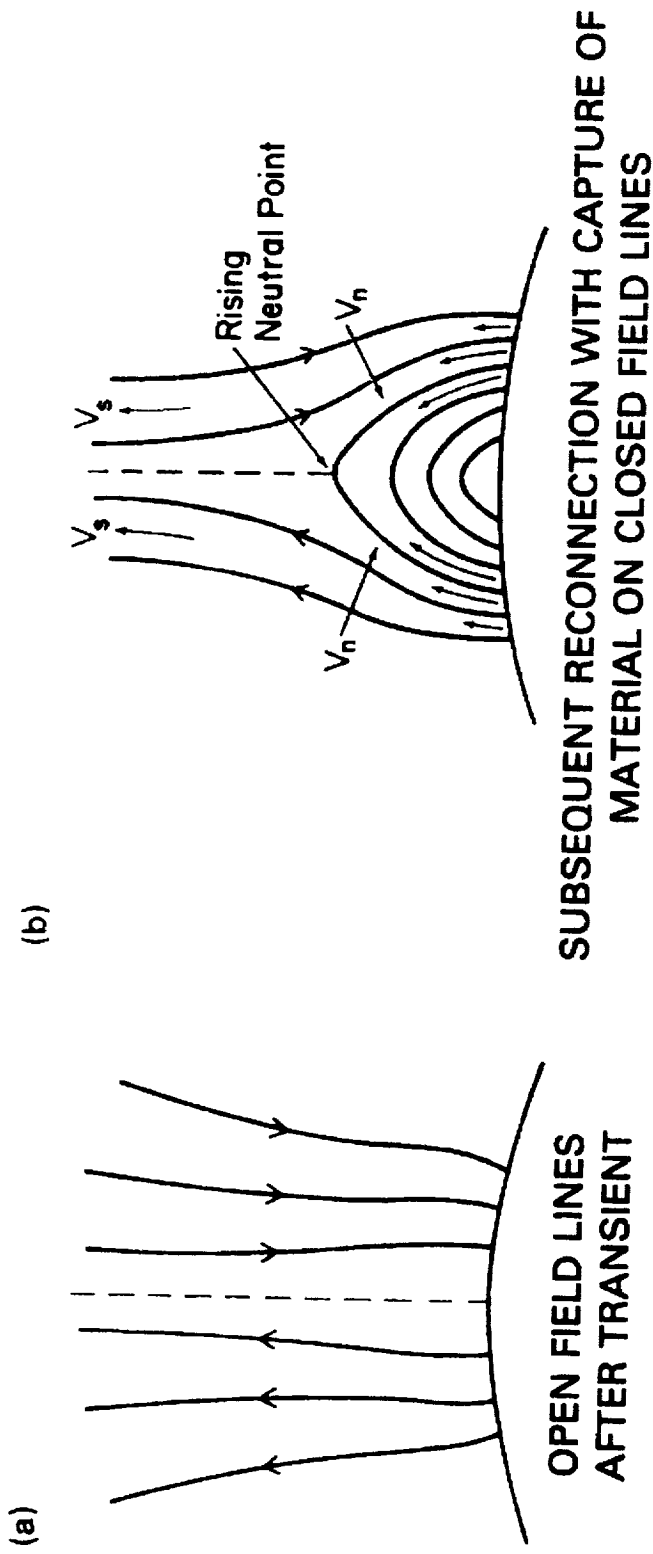


(Antiochos, DeVore, & Klimchuk 1999, *Ap. J.*, 510, 485;
Klimchuk 2001, Proc. Chapman Conf. on Space Weather, to appear)

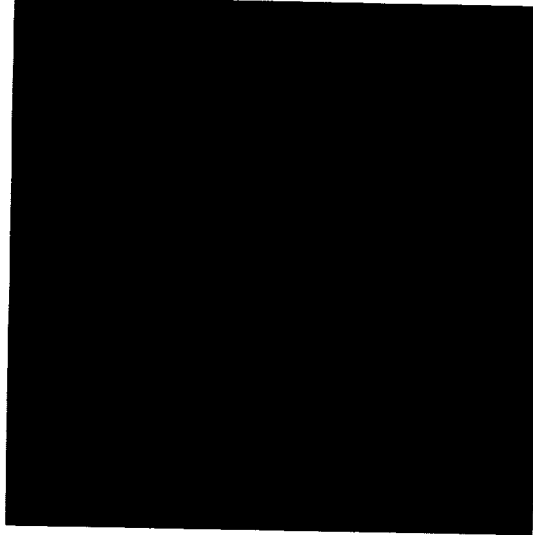


What Do We Think We Understand about CMEs?

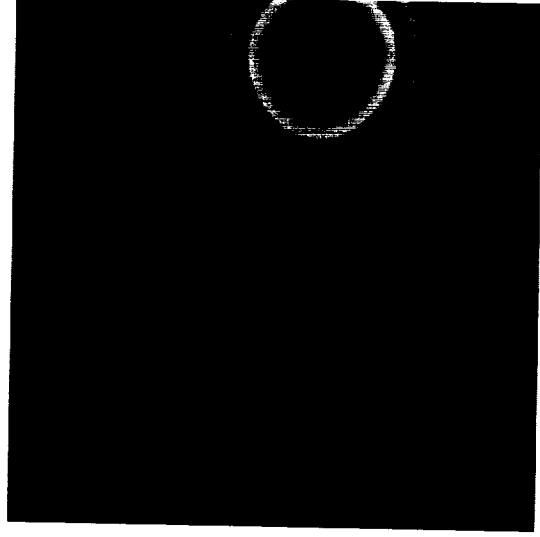
Kopp and Pneuman (1976) - "Post-Flare" Loops Formed From Reconnection



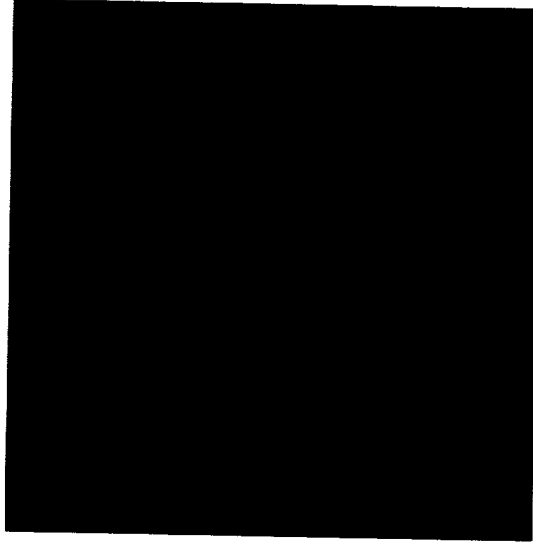
MHD Simulation of Streamer Disruption:



Equilibrium Streamer

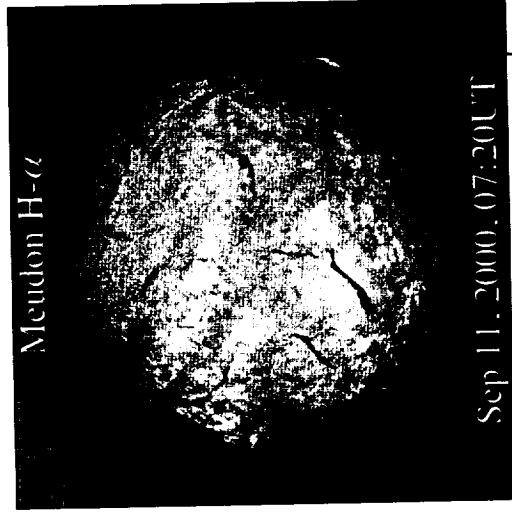
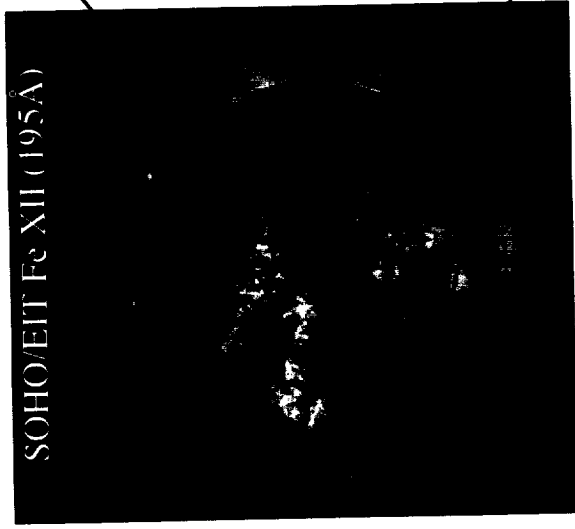


4 Hours After Eruption

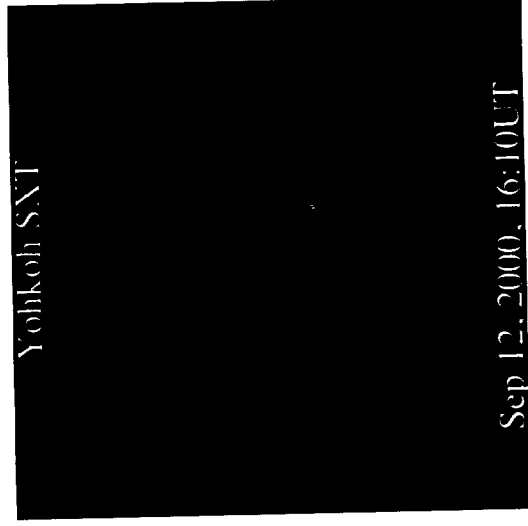


2 days After Eruption

Reconnection in the Aftermath of a CME/Prominence Eruption: Post-“flare” Loops



Pre-Eruption State (With Visible Prominence)



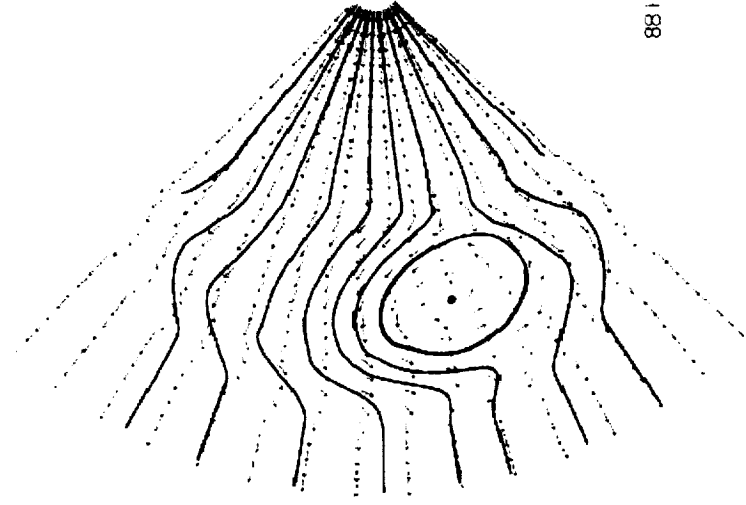
Post-Eruption State

[Play Movie ►](#)

INTERPLANETARY CONSEQUENCES

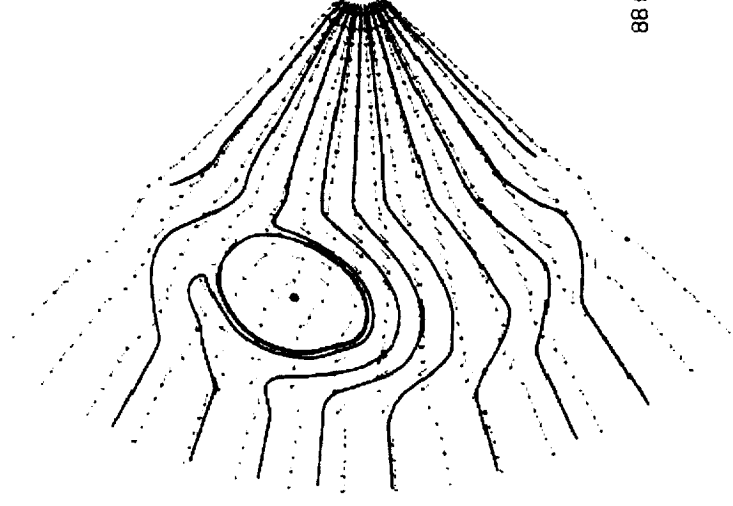
- Tremendous amount of literature on modeling flux ropes in interplanetary space (e.g., Bothmer, Burlaga, Marubashi, Osherovich, Rust)
- Computing CME evolution: It easiest to start beyond the critical points ($\geq 20 R_s$)
- Earliest work focused on interplanetary shock waves: Dryer, Wu, and co-workers
- Propagation of “spheromaks” and cylindrical flux ropes (Detman, Vandas, Odstrcil, Cargill; Recent work by Manchester et al. starting in the corona)
- How do ejecta evolve in a structured solar wind? (Odstrcil, Pizzo).
- To make the connection to eruptions seen on the Sun, we must model the CME initiation and evolution from the Sun out into the heliosphere
- Can interplanetary observations give clues to the initiation process?

From Vandas et al. (1996)



88 H

case 2a



88 H

case 2b

2D Simulation of the Propagation of Cylindrical Flux Ropes in the Solar Wind

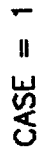
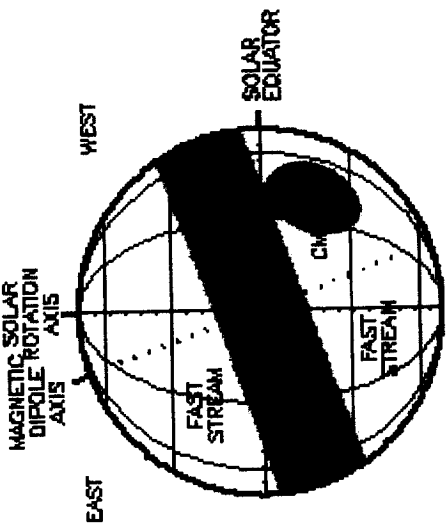
INTERPLANETARY FLUX ROPES ARISE IN COMPETING CME MODELS

- Mere presence of a flux rope is not a discriminator (different models create a flux rope prior to eruption, or in the aftermath of the eruption)
- More detailed simulations that predict more specific properties might provide discriminators (e.g., heating, composition)

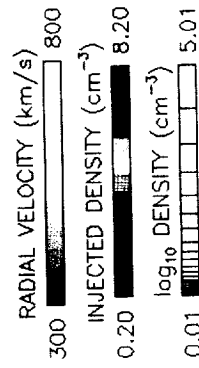
HOW WELL DO WE UNDERSTAND THE INTERPLANETARY FLUX ROPES WE SEE?

- Interplanetary flux ropes are fit quite successfully with linear force-free field models: $(\nabla \times \mathbf{B}) \times \mathbf{B} = 0$ or $(\nabla \times \mathbf{B}) = \alpha \mathbf{B}$
- $\alpha = \text{constant}$ is a major simplifying assumption
- Analyzing simulated CMEs can give us insight into the strengths and weakness of force-free models

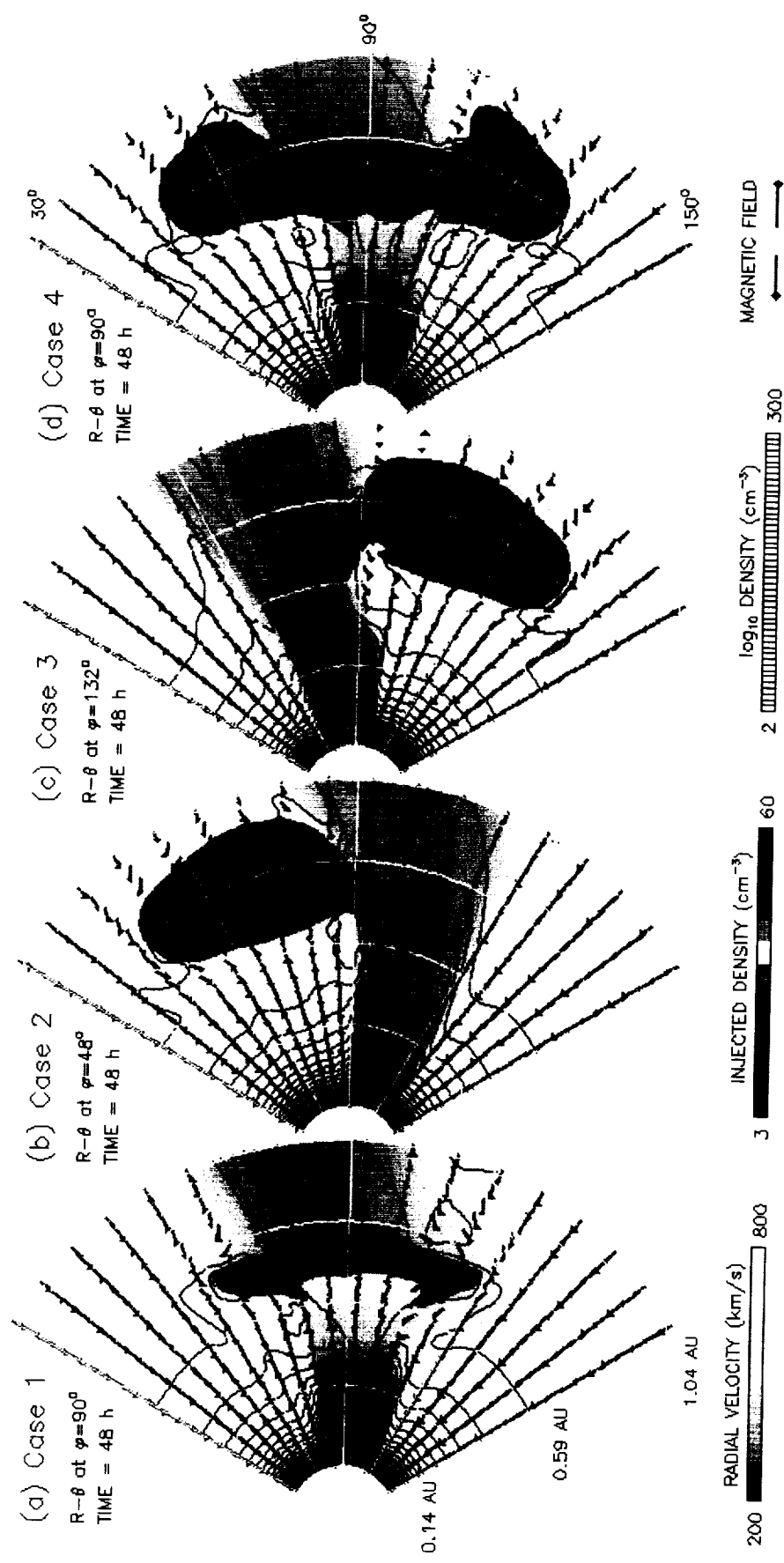
Odstrcil and Pizzo (1999ab)



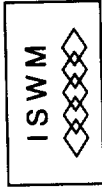
DAY = 12



From Odstreil & Pizzo (1999)

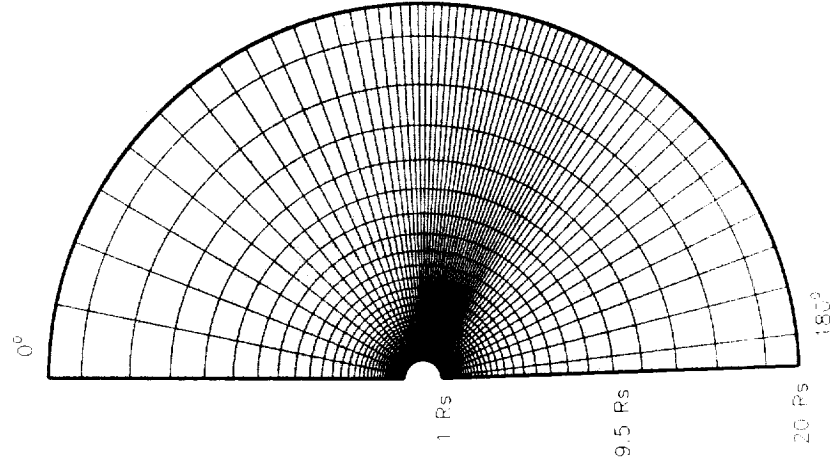


3D MHD Simulation of CME Propagation: Distortion of the Interplanetary Magnetic Field



Merged Numerical Grids

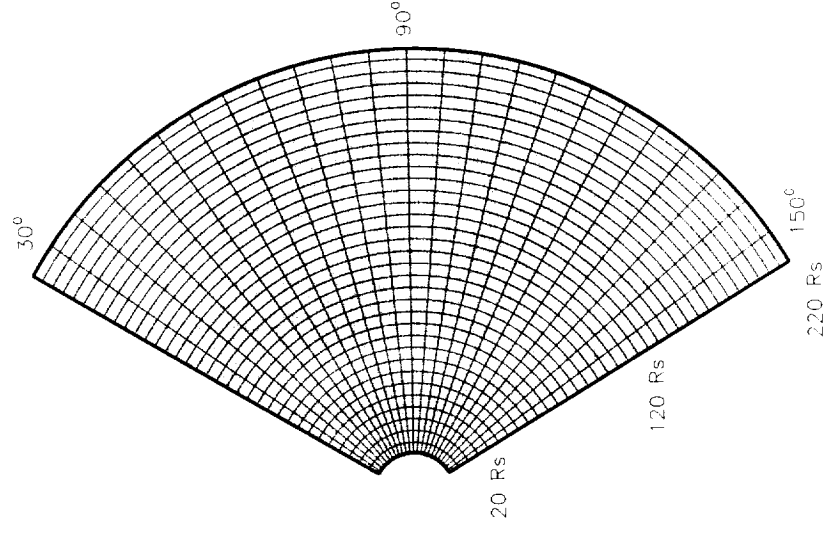
Coronal Model (1 Rs - 20 Rs)



SAIC (San Diego, CA): 200x300 grid points
 $\Delta r = 0.0053\text{--}0.59\text{ Rs}$, $\Delta\theta = 0.24\text{--}2.4^\circ$

NOTE: Only every 5th grid line is shown

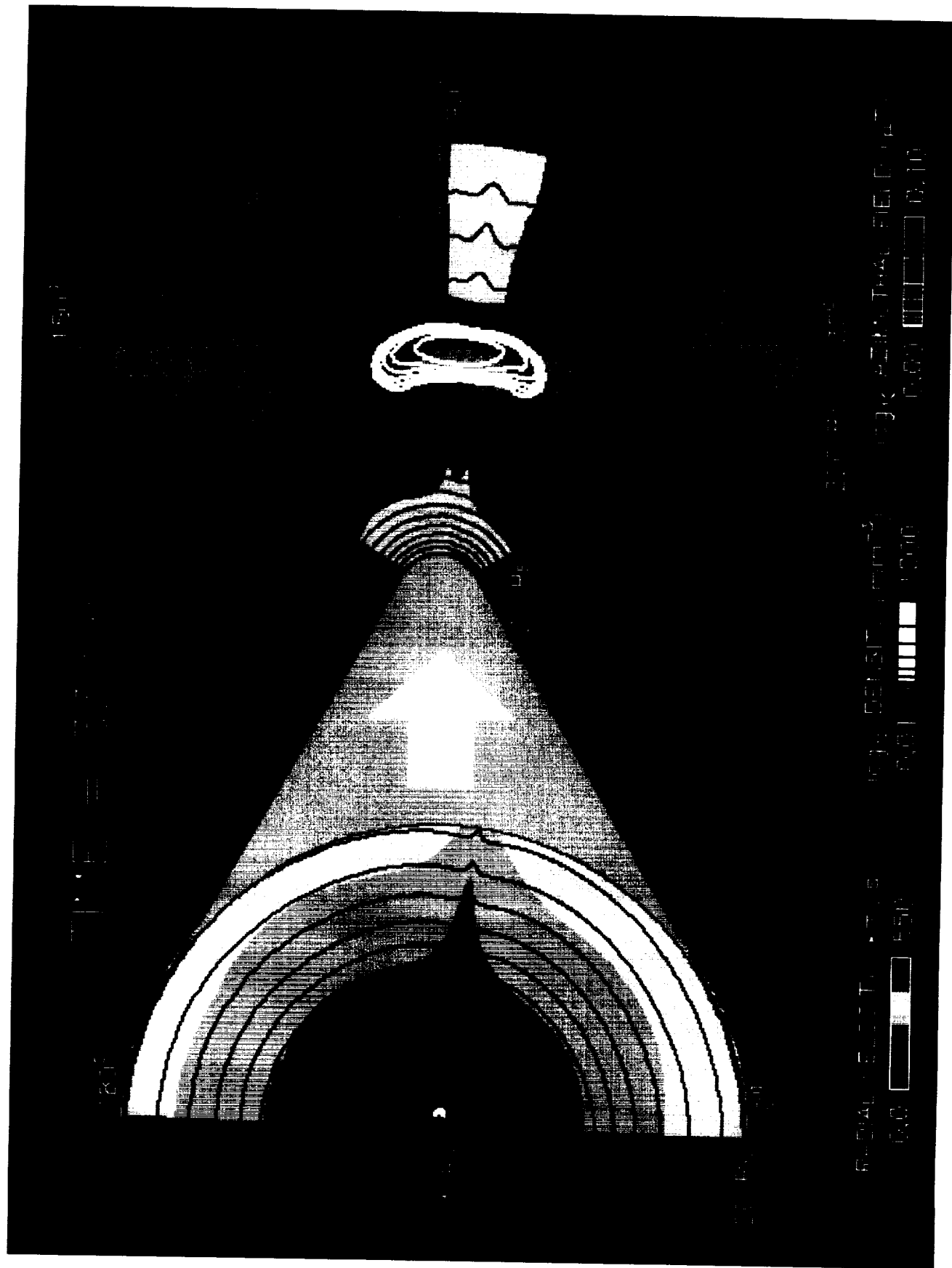
Heliospheric Model (20 Rs - 220 Rs)



CIRES/SEC (Boulder, CO): 340x240 grid points
 $\Delta r = 0.5\text{ Rs}$, $\Delta\theta = 0.5^\circ$

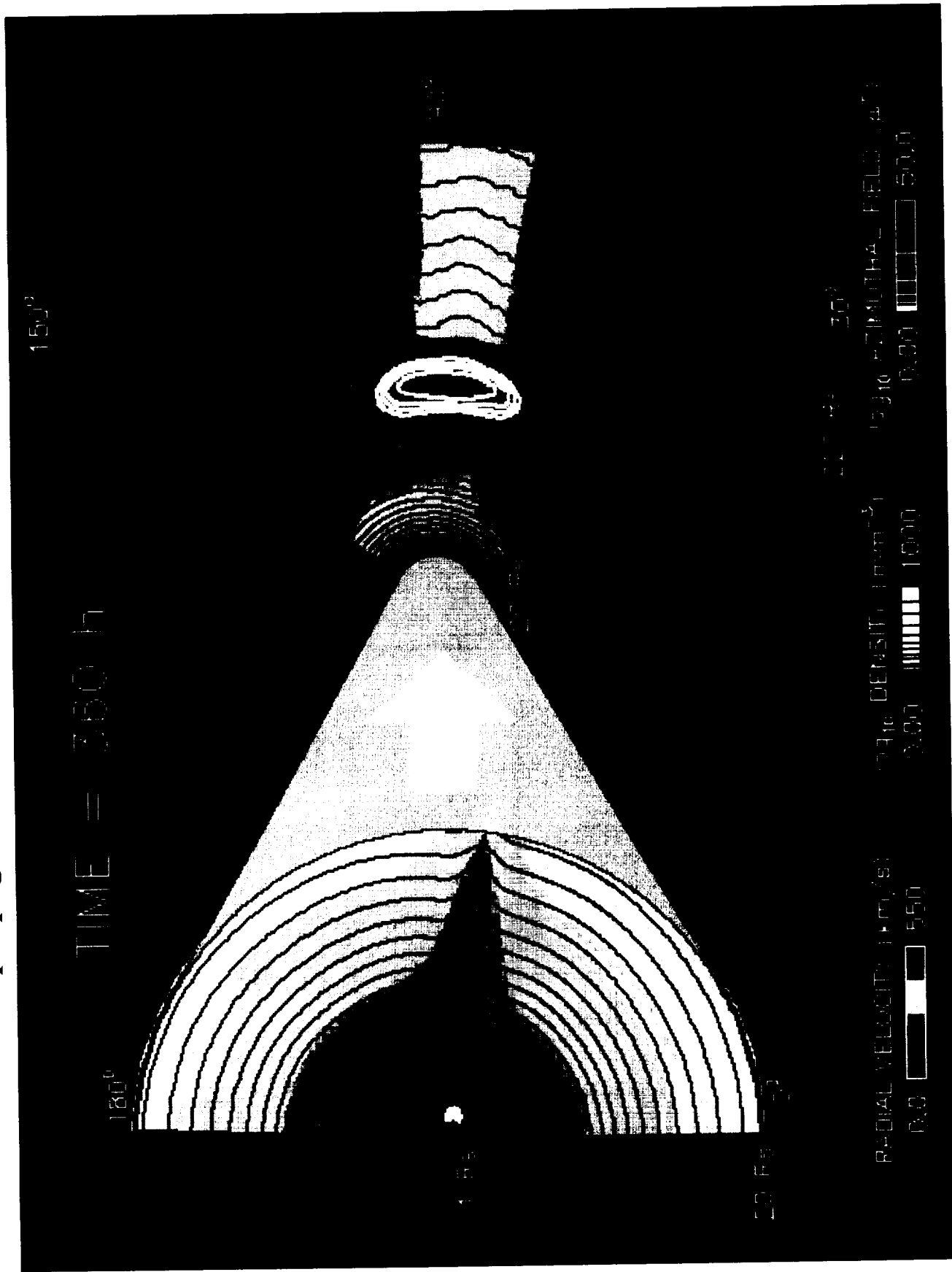
NOTE: Only every 10th grid line is shown

2D Simulation of CME propagation from the Sun to 1 A.U.: Flux Cancellation



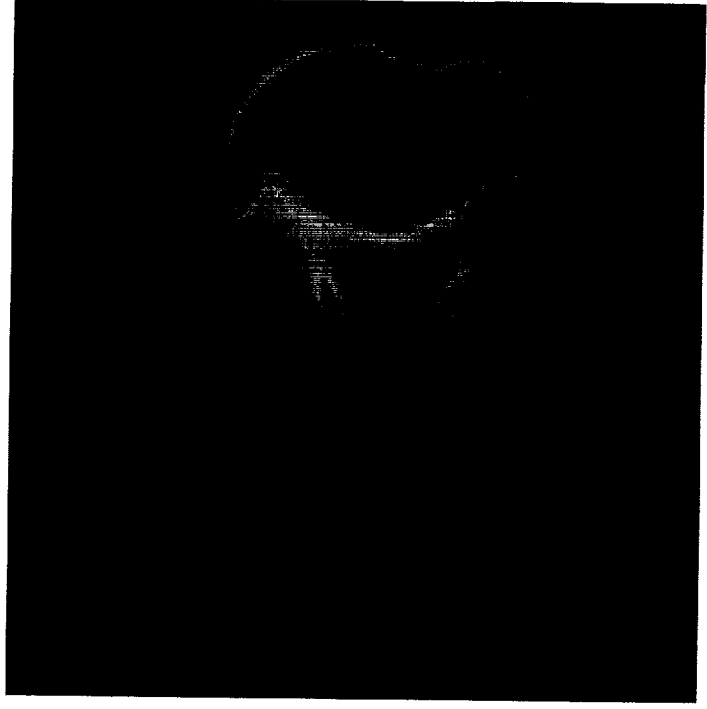
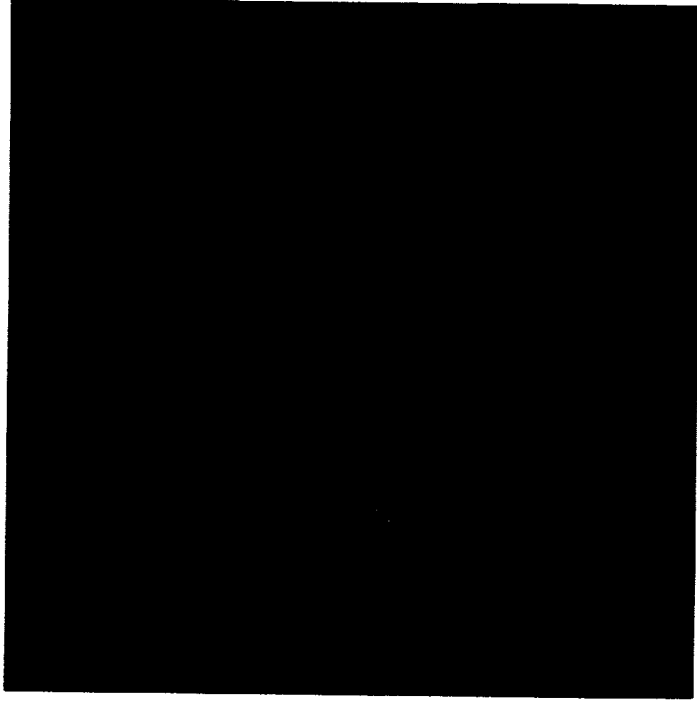
Flux Rope Forms Prior to Eruption

2D Simulation of CME propagation from the Sun to 1 A.U.: Photospheric Shearing Flows

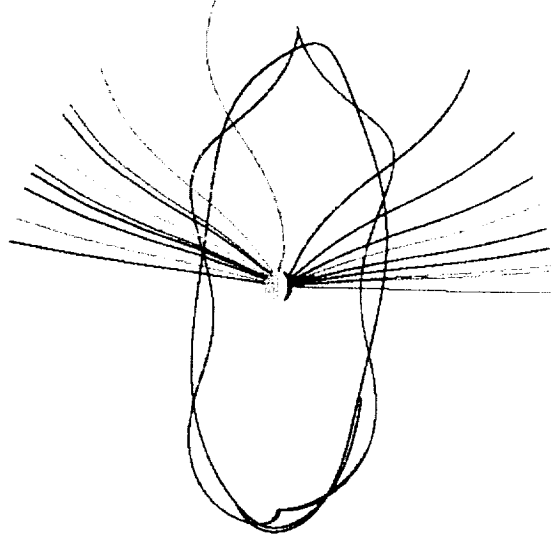
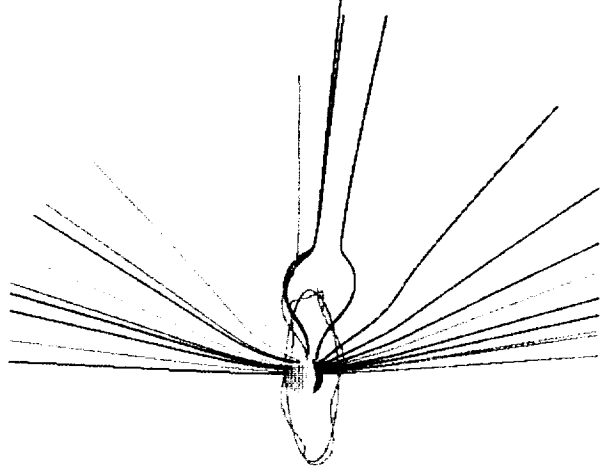
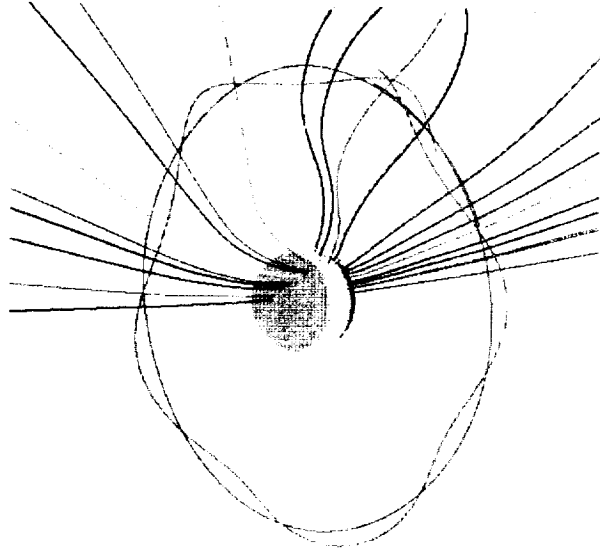


Flux Rope Forms after Eruption

CME Evolution: Near Sun

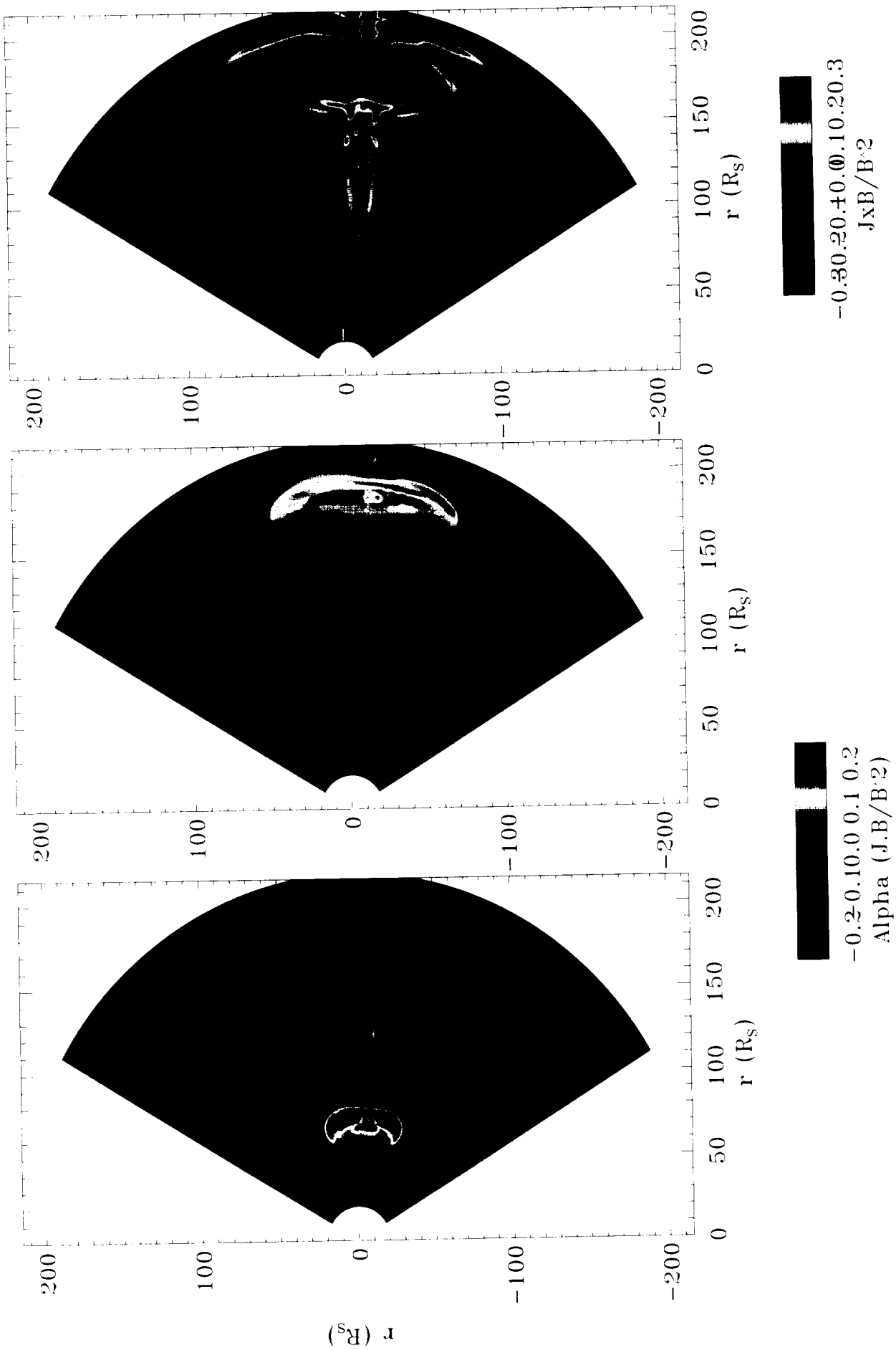


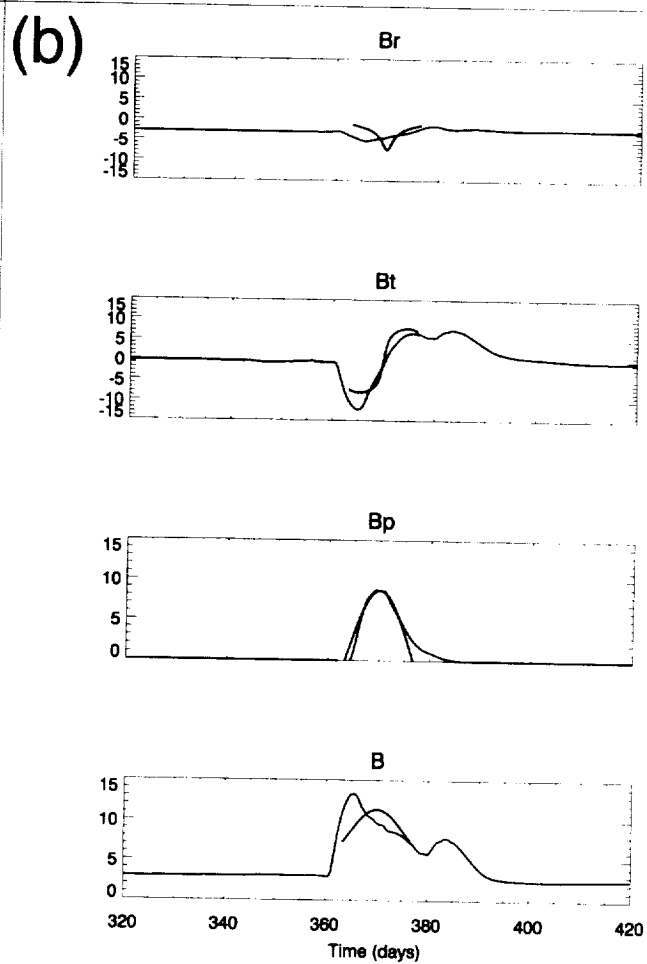
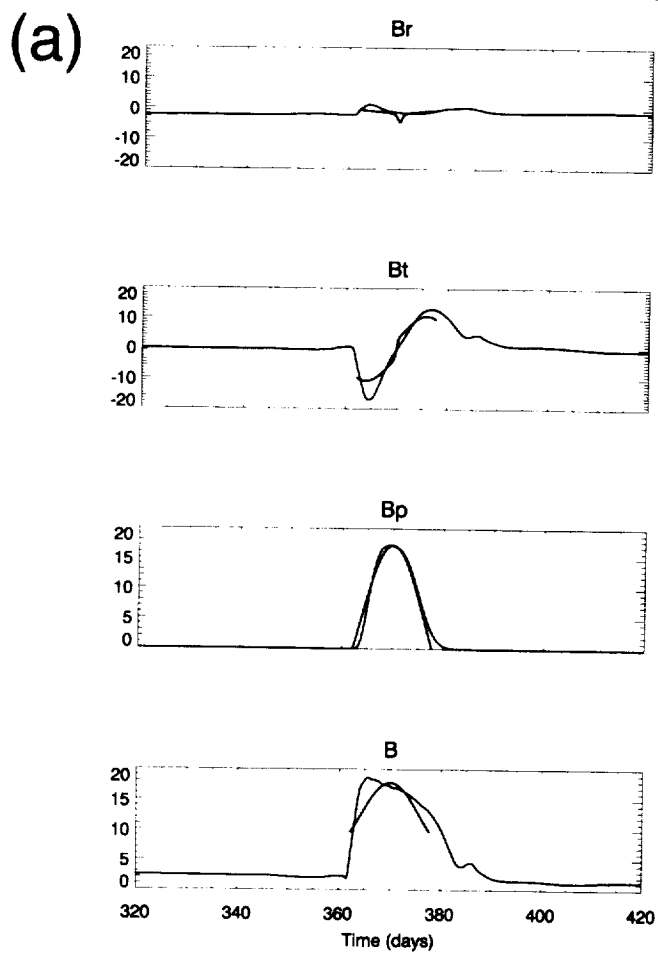
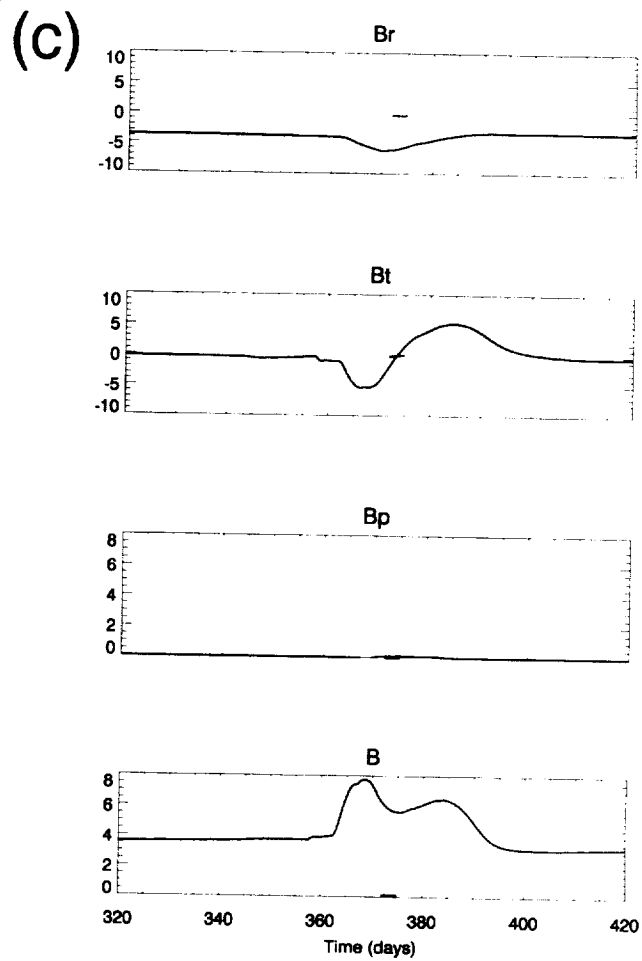
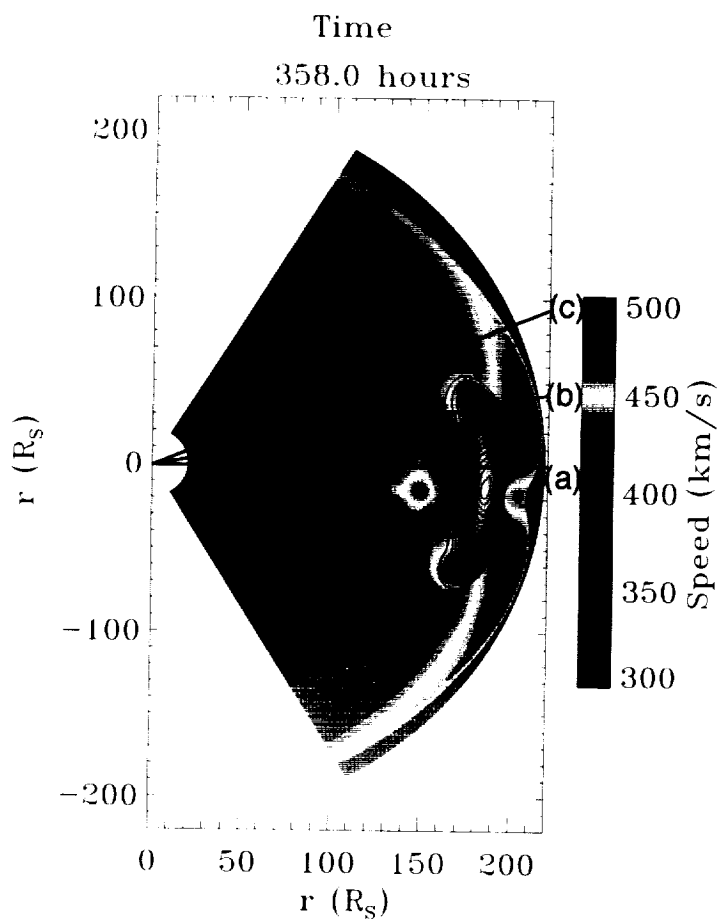
α ($= \mathbf{J} \cdot \mathbf{B} / \mathbf{B} \cdot \mathbf{B}$) evolution, scaled by αL



Flux rope field lines

α Evolution to 1 A.U.





Linear Force Free Fit Assuming Cylindrical Flux Rope:

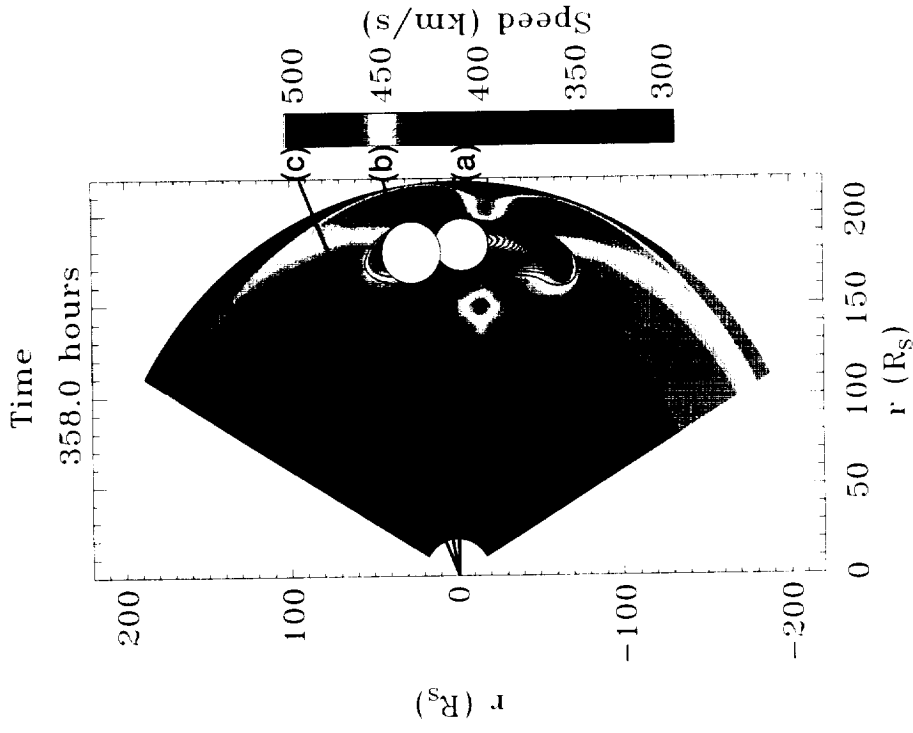
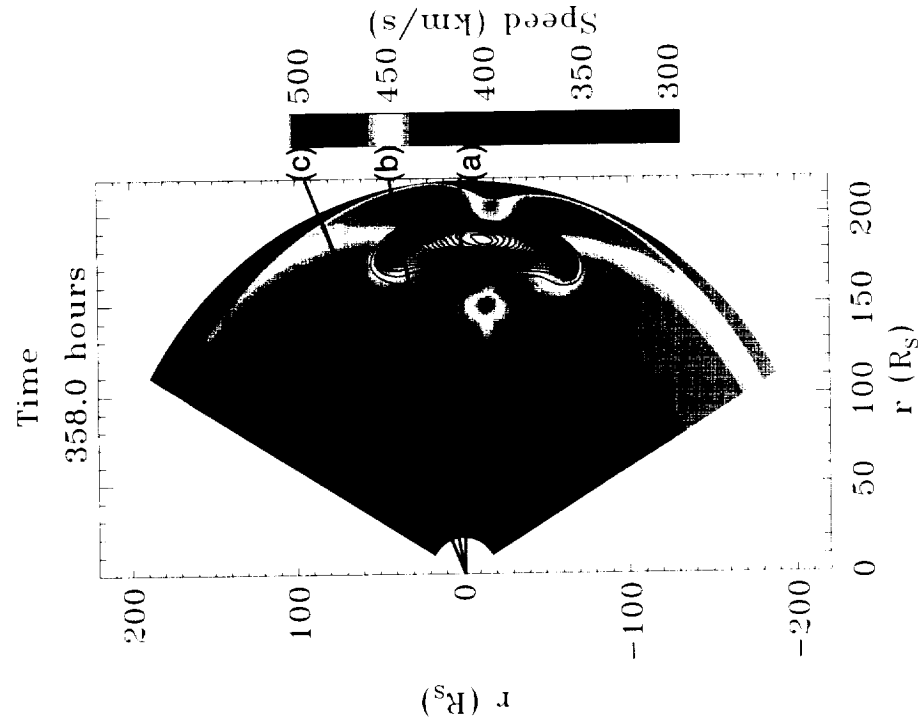
Impact parameter p :

(a) $p = .21R$, $R=30R_s$

(b) $p = .54R$, $R=36R_s$

(c) $p = 1$ (does not encounter flux rope)

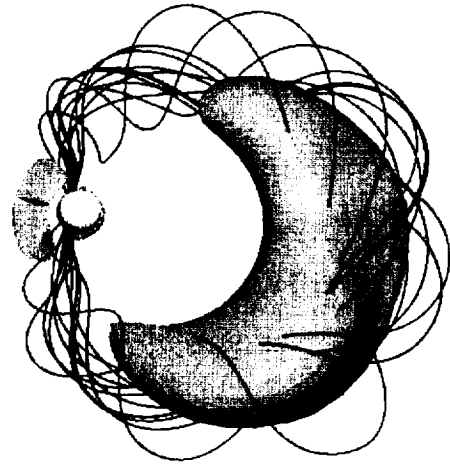
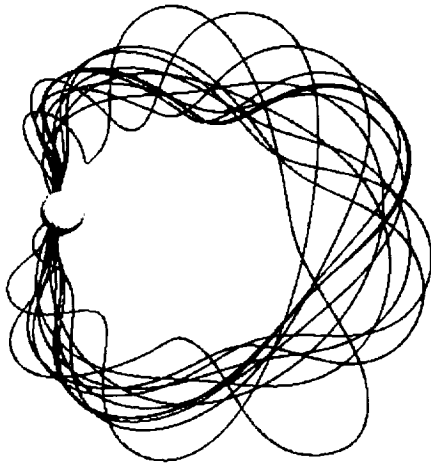
Inferred Flux Rope Shape and Size
are Misleading



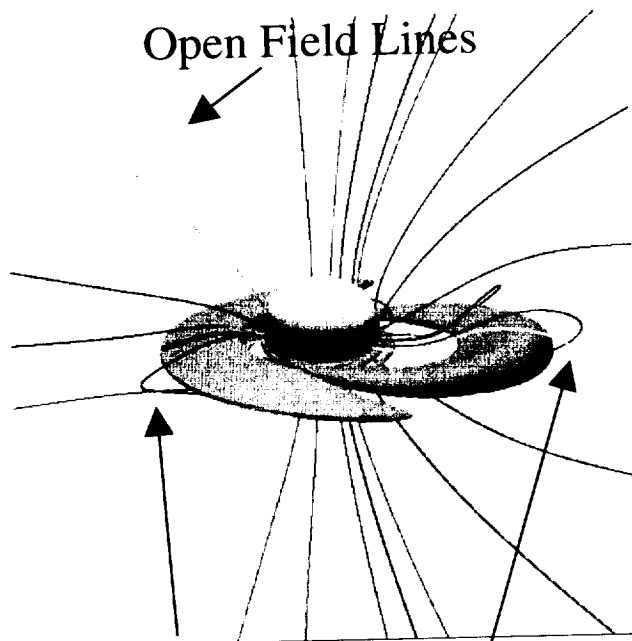
FORCE-FREE FITTING OF SIMULATED CME EJECTA

- Simulated ejecta is highly idealized, but nevertheless can yield useful insights into strengths and limitations of simple force-free fits
- Variations in force free parameter α were not too large, but α shows some non uniformity in evolution
- Force-free model fits the flux rope quite well (not surprising for 2-D: flux rope axis is known).
- Interior of flux rope is force-free, even at 1 A.U.
- Weakness of the force-free fit appears to be in the assumed shape of the flux rope
- More realistic simulations (3D, two-state wind, rotation, etc.) are required and are currently in progress.

3D CME Eruption: Magnetic Field Topology

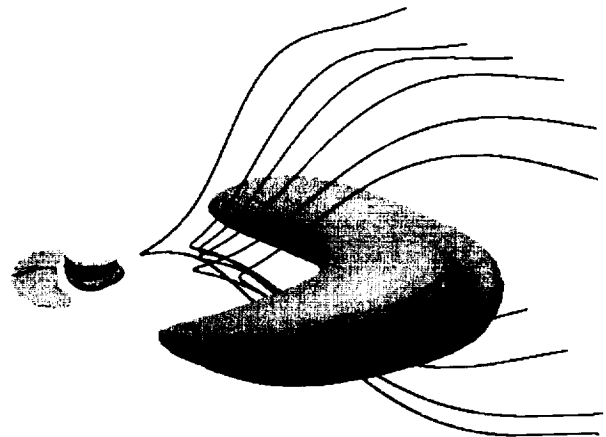


Flux Rope Connected to the Sun



Open Field Lines

Closed or Overlying Field Lines



Disconnected or
U-shaped Field Lines

THE FUTURE

- Simulation of CME propagation to 1 A.U. (and beyond) is entering a stage where real progress can be made.
- Not a moment too soon! We have many puzzles, and important upcoming observational opportunities (e.g., STEREO).
- The only way we will resolve which physical mechanism initiates CMEs will be to refine the models until they can *directly* address observations
- For example, we should try to track the evolution of an active region with detailed vector magnetograms, while comparing model output to observed quantities (e.g., X-ray emission, EUV emission). This requires significant improvements to present models.
- In situ measurements provide the ultimate test of the CME evolution predicted by the models.
- The models may in turn help us gain more insight into the interplanetary data, and devise improved analysis methods.

REPORT DOCUMENTATION PAGE

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